## Nuclear Motion in the Core-Excited States of H<sub>2</sub>O Probed by Sub-Natural-Linewidth Resonant Auger Spectroscopy

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The O 1s photoexcitation spectra of H<sub>2</sub>O exhibit two dominant peaks below threshold. The broad structureless peak at 534 eV corresponds to the O1s  $\rightarrow$  4a<sub>1</sub> excitation,whereas the peak with partially resolved vibrational structure at 536 eV is assigned to O1s  $\rightarrow$  2b<sub>2</sub>. We have investigated nuclear motions in the O 1s<sup>-1</sup>4a<sub>1</sub> and 1s<sup>-1</sup>2b<sub>2</sub> excited states using resonant Auger spectroscopy. The experiment was carried out on beamline 27SUat SPring-8.The electron spectra were collected with a high-resolution electron spectrometer SES-2002 (Gammadata-Scienta) under sub-natural-linewidth conditions, i.e.with overall width (< 100 meV) smaller than the lifetime broadening of the core-excited states (~150 meV).

.The resonance enhancement of photoemission intensity from the  $1b_1$  and  $1b_2$  orbitals are significant. The progression of the symmetric stretching vibrations is clearly seen in the  $1b_1$  photoemission band, whereas the vibrational members are only partially resolved in the  $1b_2$  photoemission band because of overlap between bending and stretching modes. It is interesting to notice that the relative intensity of the symmetric stretching mode versus the bending mode changes as a function of photon energy, .e. when different portions of the potential surface of the intermediate state are sampled. Note also that the population of high vibrational members increases with the increase in the vibrational energy in the core-excited state. The reason why vibrations are highly excited in the Auger final state is that the coreexcited state is unstable at the stable point of the ground state along the stretching and bending coordinates. Therefore nuclear motions along these coordinates proceed in the coreexcited state before they are transferred to the Auger final state.

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